

On Airy Solutions of the Second Painlevé Equation

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Abstract

In this paper we discuss Airy solutions of the second Painlevé equation (P_{II}) and two related equations, the Painlevé XXXIV equation (P_{34}) and the Jimbo-Miwa-Okamoto σ form of P_{II} (S_{II}), are discussed. It is shown that solutions which depend only on the Airy function $Ai(z)$ have a completely difference structure to those which involve a linear combination of the Airy functions $Ai(z)$ and $Bi(z)$. For all three equations, the special solutions which depend only on $Ai(t)$ are *tronquée* solutions, i.e. they have no poles in a sector of the complex plane. Further for both P_{34} and S_{II} , it is shown that amongst these *tronquée* solutions there is a family of solutions which have no poles on the real axis.

Dedicated to Mark Ablowitz on his 70th birthday

1 Introduction

The six Painlevé equations (P_I – P_{VI}) were first discovered by Painlevé, Gambier and their colleagues in an investigation of which second order ordinary differential equations of the form

$$\frac{d^2q}{dz^2} = F\left(\frac{dq}{dz}, q, z\right), \quad (1.1)$$

where F is rational in dq/dz and q and analytic in z , have the property that their solutions have no movable branch points. They showed that there were fifty canonical equations of the form (1.1) with this property, now known as the *Painlevé property*. Further Painlevé, Gambier and their colleagues showed that of these fifty equations, forty-four can be reduced to linear equations, solved in terms of elliptic functions, or are reducible to one of six new nonlinear ordinary differential equations that define new transcendental functions, see Ince [1]. The Painlevé equations can be thought of as nonlinear analogues of the classical special functions [2, 3, 4, 5, 6], and arise in a wide variety of applications, for example random matrices, cf. [7, 8].

In this paper we are concerned with special solutions of the second Painlevé equation (P_{II})

$$\frac{d^2q}{dz^2} = 2q^3 + zq + \alpha, \quad (1.2)$$

with α an arbitrary constant, and two related equations. These are the Painlevé XXXIV equation (P_{34})

$$\frac{d^2p}{dz^2} = \frac{1}{2p} \left(\frac{dp}{dz}\right)^2 + 2p^2 - zp - \frac{(\alpha + \frac{1}{2})^2}{2p}, \quad (1.3)$$

is equivalent to equation XXXIV of Chapter 14 in [1], which is solvable in terms of P_{II} (1.2), see §2.1, and the Jimbo-Miwa-Okamoto σ form of P_{II} (S_{II})

$$\left(\frac{d^2\sigma}{dz^2}\right)^2 + 4\left(\frac{d\sigma}{dz}\right)^3 + 2\frac{d\sigma}{dz} \left(z\frac{d\sigma}{dz} - \sigma\right) = \frac{1}{4}(\alpha + \frac{1}{2})^2, \quad (1.4)$$

which is satisfied by the Hamiltonian associated with P_{II} (1.2), see §2.1. Equation (1.4) is equation SD-I.d in the classification of second-order, second-degree equations which have the Painlevé property by Cosgrove and Scoufis

[9], an equation first derived by Chazy [10]. Frequently in applications it is the associated second-order, second-degree equation such as (1.4) which arises rather than the Painlevé equation.

It is well-known that P_{II} has special solutions depending on one parameter that are expressed in terms of the Airy functions $Ai(z)$ and $Bi(z)$. In this paper we study the Airy solutions for P_{II} (1.2), P_{34} (1.3) and S_{II} (1.4), see §3. In particular it is shown that the solutions which depend only on $Ai(z)$ have a completely different asymptotic behaviour as $z \rightarrow -\infty$ in comparison to those which involve a linear combination of $Ai(z)$ and $Bi(z)$, which is a new characterization of these special solutions of P_{II} (1.2), P_{34} (1.3) and S_{II} (1.4). Further, the special solutions of the three equations which depend only on $Ai(t)$ are *tronquée* solutions, i.e. they have no poles in a sector of the complex plane. Additionally it is shown that there are new families of Airy solutions of P_{34} (1.3) and S_{II} (1.4) which have no poles on the real axis and so are likely to arise in physical applications. For P_{34} (1.3) these solutions have algebraic decay as $z \rightarrow \pm\infty$, whilst for S_{II} (1.4) these solutions have algebraic decay as $z \rightarrow +\infty$ and algebraic growth as $z \rightarrow -\infty$.

2 Some properties of the second Painlevé equation

2.1 Hamiltonian structure

Each of the Painlevé equations P_I – P_{VI} can be written as a Hamiltonian system

$$\frac{dq}{dz} = \frac{\partial \mathcal{H}_J}{\partial p}, \quad \frac{dp}{dz} = -\frac{\partial \mathcal{H}_J}{\partial q}, \quad (2.1)$$

for a suitable Hamiltonian function $\mathcal{H}_J(q, p, z)$ [11, 12, 13, 14]. The function $\sigma(z) \equiv \mathcal{H}_J(q, p, z)$ satisfies a second-order, second-degree ordinary differential equation, whose solution is expressible in terms of the solution of the associated Painlevé equation [11, 12, 13, 14].

The Hamiltonian associated with P_{II} (1.2) is

$$\mathcal{H}_{II}(q, p, z; \alpha) = \frac{1}{2}p^2 - (q^2 + \frac{1}{2}z)p - (\alpha + \frac{1}{2})q \quad (2.2)$$

and so

$$\frac{dq}{dz} = p - q^2 - \frac{1}{2}z, \quad \frac{dp}{dz} = 2qp + \alpha + \frac{1}{2} \quad (2.3)$$

(Jimbo and Miwa [11], Okamoto [14]). Eliminating p in (2.3) then q satisfies P_{II} (1.2) whilst eliminating q yields (1.3). Further if q satisfies P_{II} (1.2) then $p = q' + q^2 + \frac{1}{2}z$ satisfies (1.3). Conversely if p satisfies (1.3) then $q = (p' - \alpha - \frac{1}{2})/(2p)$ satisfies P_{II} (1.2). Thus there is a one-to-one correspondence between solutions of P_{II} (1.2) and those of P_{34} (1.3).

An important property of the Hamiltonian (2.2), which is very useful in applications, is that it satisfies a second-order, second-degree ordinary differential equation, as discussed in the following theorem.

Theorem 2.1. *Consider the function $\sigma(z; \alpha) = \mathcal{H}_{II}(q, p, z; \alpha)$ defined by (2.2), where q and p satisfy the system (2.3), then $\sigma(z; \alpha)$ satisfies (1.4). Conversely if $\sigma(z; \alpha)$ is a solution of (1.4), then*

$$q(z; \alpha) = \frac{4\sigma''(z; \alpha) + 2\alpha + 1}{8\sigma'(z; \alpha)}, \quad p(z; \alpha) = -2\sigma'(z; \alpha), \quad (2.4)$$

with $' \equiv d/dz$, are solutions of (1.2) and (1.3), respectively.

Proof. See Jimbo and Miwa [11] and Okamoto [12, 13, 14]. □

2.2 Bäcklund transformations

The Painlevé equations P_{II} – P_{VI} possess *Bäcklund transformations* which relate one solution to another solution either of the same equation, with different values of the parameters, or another equation (see [2, 4, 15, 16] and the references therein). An important application of the Bäcklund transformations is that they generate hierarchies of classical solutions of the Painlevé equations, which are discussed in §3.

The Bäcklund transformations for P_{II} (1.2) are given in the following theorem.

Theorem 2.2. Let $q \equiv q(z; \alpha)$ is a solution of P_{II} (1.2), then the transformations

$$\mathcal{S} : \quad q(z; -\alpha) = -q, \quad (2.5)$$

$$\mathcal{T}_{\pm} : \quad q(z; \alpha \pm 1) = -q - \frac{2\alpha \pm 1}{2q^2 \pm 2q' + z}, \quad (2.6)$$

give solutions of P_{II} , provided that $\alpha \neq \mp \frac{1}{2}$ in (2.6).

Proof. See Gambier [17]. □

The solutions $q_{\alpha} = q(z; \alpha)$, $q_{\alpha \pm 1} = q(z; \alpha \pm 1)$ also satisfy the nonlinear recurrence relation

$$\frac{2\alpha + 1}{q_{\alpha+1} + q_{\alpha}} + \frac{2\alpha - 1}{q_{\alpha} + q_{\alpha-1}} + 4q_{\alpha}^2 + 2z = 0, \quad (2.7)$$

a difference equation which is known as an alternative form of discrete P_I (alt-d P_I) [18]. The difference equation (2.7) is obtained by eliminating q' between the transformations \mathcal{T}_{\pm} given by (2.6). Note that for P_{II} (1.2), the independent variable z varies and the parameter α is fixed, whilst for the discrete equation (2.7), z is a fixed parameter and α varies.

3 Special function solutions

The Painlevé equations P_{II} – P_{VI} possess hierarchies of solutions expressible in terms of classical special functions, for special values of the parameters through an associated Riccati equation,

$$\frac{dq}{dz} = f_2(z)q^2 + f_1(z)q + f_0(z), \quad (3.1)$$

where $f_2(z)$, $f_1(z)$ and $f_0(z)$ are rational functions. Hierarchies of solutions, which are often referred to as “one-parameter solutions” (since they have one arbitrary constant), are generated from “seed solutions” derived from the Riccati equation using the Bäcklund transformations given in §2.2. Furthermore, as for the rational solutions, these special function solutions are often expressed in the form of determinants.

Solutions of P_{II} – P_{VI} are expressed in terms of special functions as follows (see [2, 4, 15, 19], and the references therein): for P_{II} in terms of Airy functions $\text{Ai}(z)$, $\text{Bi}(z)$; for P_{III} in terms of Bessel functions $J_{\nu}(z)$, $Y_{\nu}(z)$; for P_{IV} in terms of parabolic cylinder functions $D_{\nu}(z)$; for P_V in terms of confluent hypergeometric functions ${}_1F_1(a; c; z)$, equivalently Kummer functions $M(a, b, z)$, $U(a, b, z)$ or Whittaker functions $M_{\kappa, \mu}(z)$, $q_{\kappa, \mu}(z)$; and for P_{VI} in terms of hypergeometric functions ${}_2F_1(a, b; c; z)$. Some classical orthogonal polynomials arise as particular cases of these special functions and thus yield rational solutions of the associated Painlevé equations: for P_{III} and P_V in terms of associated Laguerre polynomials $L_k^{(m)}(z)$; for P_{IV} in terms of Hermite polynomials $H_n(z)$; and for P_{VI} in terms of Jacobi polynomials $P_n^{(\alpha, \beta)}(z)$.

3.1 Second Painlevé equation

We note that P_{II} (1.2) can be written as

$$\varepsilon \frac{d}{dz} \left(\varepsilon \frac{dq}{dz} - q^2 - \frac{1}{2}z \right) + 2q \left(\varepsilon \frac{dq}{dz} - q^2 - \frac{1}{2}z \right) = \alpha - \frac{1}{2}\varepsilon,$$

with $\varepsilon^2 = 1$. Hence if $\alpha = \frac{1}{2}\varepsilon$, then special solutions of P_{II} can be obtained in terms of solutions of the Riccati equation

$$\varepsilon \frac{dq}{dz} = q^2 + \frac{1}{2}z. \quad (3.2)$$

Any solution of this equation is also a solution of P_{II} (1.2), provided that $\alpha = \frac{1}{2}\varepsilon$. Linearising the Riccati equation (3.2) by setting $q = -\varepsilon\varphi'/\varphi$ yields

$$\frac{d^2\varphi}{dz^2} + \frac{1}{2}z\varphi = 0, \quad (3.3)$$

which is equivalent to the Airy equation and has general solution

$$\varphi(z; \vartheta) = \cos(\vartheta) \text{Ai}(t) + \sin(\vartheta) \text{Bi}(t), \quad t = -2^{-1/3}z, \quad (3.4)$$

with $\text{Ai}(t)$ and $\text{Bi}(t)$ the Airy functions and ϑ an arbitrary constant. The Airy solutions of P_{II} are classified in the following theorem due to Gambier [17].

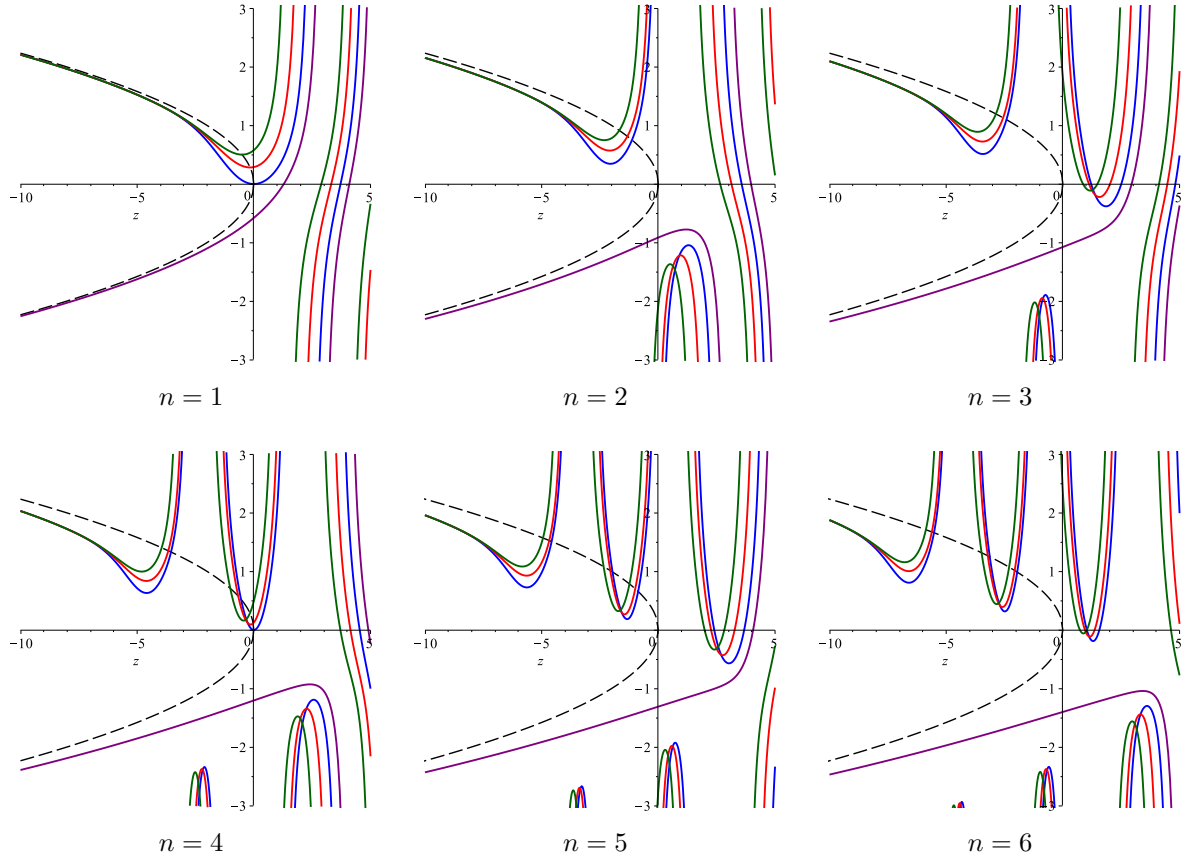


Figure 3.1: Plots of $q_n(z; \vartheta)$ (3.7) for $\vartheta = 0$ [purple], $\vartheta = \frac{1}{6}\pi$ [blue], $\vartheta = \frac{1}{3}\pi$ [red], $\vartheta = \frac{1}{2}\pi$ [green]; the dashed line is the parabola $2q^2 + z = 0$.

Theorem 3.1. *The second Painlevé equation (1.2) has a one-parameter family of solutions expressible in terms of Airy functions given by (3.4) if and only if $\alpha = n - \frac{1}{2}$, with $n \in \mathbb{Z}$.*

The simplest Airy solutions of P_{II} (1.2), which arise when $\alpha = \pm \frac{1}{2}$, are

$$q(z) = \mp \frac{d}{dz} \ln \varphi(z; \vartheta), \quad (3.5)$$

with $\varphi(z; \vartheta)$ given by (3.4). Then using the Bäcklund transformations (2.6), a hierarchy of Airy solutions for $\alpha = n + \frac{1}{2}$, $n \in \mathbb{Z}$, can be generated. The Airy solutions can also be expressed in terms of determinants, as described in the following theorem.

Theorem 3.2. *Let $\tau_n(z; \vartheta)$ be the $n \times n$ determinant*

$$\tau_n(z; \vartheta) = \left[\frac{d^{j+k}}{dz^{j+k}} \varphi(z; \vartheta) \right]_{j,k=0}^{n-1}, \quad n \geq 1, \quad (3.6)$$

with $\varphi(z; \vartheta)$ given by (3.4) and $\tau_0(z; \vartheta) = 1$, then

$$q_n(z; \vartheta) = \frac{d}{dz} \ln \frac{\tau_{n-1}(z; \vartheta)}{\tau_n(z; \vartheta)}, \quad n \geq 1, \quad (3.7)$$

satisfies P_{II} (1.2) with $\alpha = n - \frac{1}{2}$.

Proof. See Flaschka and Newell [20], Okamoto [14]; also [21]. □

n	$q_n(z; \vartheta)$
1	Φ
2	$-\Phi - \frac{1}{2\Phi^2 + z}$
3	$\frac{1}{2\Phi^2 + z} - \frac{2z\Phi^2 - \Phi + z^2}{4\Phi^3 + 2z\Phi + 1}$
4	$\frac{2z\Phi^2 - \Phi + z^2}{4\Phi^3 + 2z\Phi + 1} - \frac{48\Phi^3 - 8z^2\Phi^2 + 28z\Phi - 4z^3 + 9}{z(8z\Phi^4 - 16\Phi^3 + 8z^2\Phi^2 - 8z\Phi + 2z^3 - 3)} - \frac{3}{z}$

Table 3.1: The Airy solutions $q_n(z; \vartheta)$ of P_{II} (1.2), see (3.7), where $\Phi = -\varphi'(z; \vartheta)/\varphi(z; \vartheta)$, with $\varphi(z; \vartheta)$ given by (3.4).

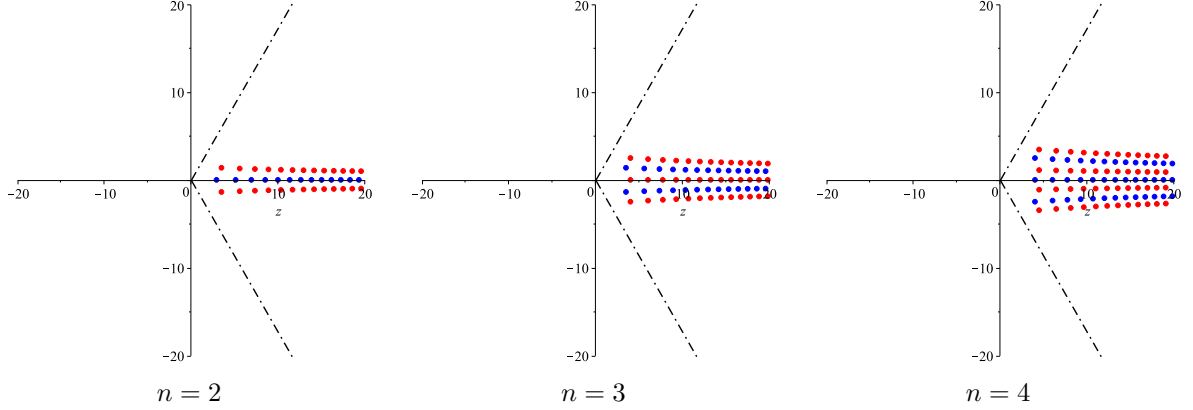


Figure 3.2: Plots of the poles of $q_n(z; 0)$ and $\sigma_n(z; 0)$ for $n = 2, 3, 4$; the blue and red circles represent poles with residues $+1$ and -1 , respectively.

We remark that the determinant

$$\Delta_n(t) = \det \left[\frac{d^{j+k}}{dt^{j+k}} \text{Ai}(t) \right]_{j,k=0}^{n-1},$$

which is equivalent to $\tau_n(z; 0)$ given by (3.6), arises in random matrix theory, in connection with the Gaussian Unitary Ensemble (GUE) in the soft-edge scaling limit, see e.g. [21, p. 393].

If we set $\Phi(z; \vartheta) \equiv -\varphi'(z; \vartheta)/\varphi(z; \vartheta)$, with $\varphi(z)$ given by (3.4), then the first few solutions in the Airy function solution hierarchy for P_{II} (1.2) are given in Table 3.1. We note that $\Phi(z; \vartheta)$ satisfies the Riccati equation (3.2) with $\varepsilon = 1$.

Plots of the solutions $q_n(z; \vartheta)$ (3.7) for various ϑ are given in Figure 3.1. These plots show that the asymptotic behaviour as $z \rightarrow -\infty$ of the solutions is completely different in the case when $\vartheta = 0$ compared to the case when $\vartheta \neq 0$, see Theorem 3.3. Fornberg and Weideman [22, Figure 3] plot the locations of the poles for the solution $q_3(z; \vartheta)$ of P_{II} (1.2), i.e. when $\alpha = \frac{5}{2}$, for various choices of ϑ , which show that the pole structure of the solutions is significantly different in the case when $\vartheta = 0$ compared to the case when $\vartheta \neq 0$. The solution $q_3(z; 0)$, which depends only on $\text{Ai}(t)$, is a *tronquée* solution, i.e. it has no poles in a sector of the complex plane; such solutions play an important role in the theory of Painlevé equations. Numerical calculations suggest that the special solutions $q_n(z; 0)$ are *tronquée* solutions, see Figure 3.2.

When $\vartheta = 0$ the solution $q_n(z; 0)$ involves only on the Airy function $\text{Ai}(t)$, with $t = -2^{-1/3}z$, whereas the solution $q_n(z; \vartheta)$, with $\vartheta \neq 0$, involves the Airy function $\text{Bi}(t)$. The known asymptotics of $\text{Ai}(t)$ and $\text{Bi}(t)$ as $t \rightarrow \infty$

n	$p_n(z; \vartheta)$
1	$2\Phi^2 + z$
2	$\frac{4\Phi}{2\Phi^2 + z} + \frac{2}{(2\Phi^2 + z)^2}$
3	$z - \frac{8\Phi^2 - 2z^2\Phi + 6z}{4\Phi^3 + 2z\Phi + 1} + \frac{2(2z^2 + 1)\Phi^2 - 2z^2\Phi + 2z(z^3 - 1)}{(4\Phi^3 + 2z\Phi + 1)^2}$

Table 3.2: The Airy solutions $p_n(z; \vartheta)$ of P_{34} (1.3) given by (3.9).

are

$$\text{Ai}(t) = \frac{1}{2}\pi^{-1/2} t^{-1/4} \exp\left(-\frac{2}{3}t^{3/2}\right) \left\{1 + \mathcal{O}(t^{-3/2})\right\}, \quad (3.8a)$$

$$\text{Bi}(t) = \pi^{-1/2} t^{-1/4} \exp\left(\frac{2}{3}t^{3/2}\right) \left\{1 + \mathcal{O}(t^{-3/2})\right\}, \quad (3.8b)$$

cf. [23, §9.7(ii)]. Consequently the asymptotic behaviour of $q_n(z; \vartheta)$ as $z \rightarrow -\infty$ critically depends on whether it involves $\text{Bi}(t)$.

The asymptotic behaviour of the Airy solutions $q_n(z; \vartheta)$ as $z \rightarrow -\infty$ is given in the following theorem.

Theorem 3.3. *Let $q_n(z; \vartheta)$ be defined by (3.7), then as $z \rightarrow -\infty$,*

$$q_n(z; \vartheta) = \begin{cases} -\frac{(-z)^{1/2}}{\sqrt{2}} + \frac{2n-1}{4z} + \frac{12n^2 - 12n + 5}{16\sqrt{2}(-z)^{5/2}} + \mathcal{O}(z^{-4}), & \text{if } \vartheta = 0, \\ \frac{(-z)^{1/2}}{\sqrt{2}} + \frac{2n-1}{4z} - \frac{12n^2 - 12n + 5}{16\sqrt{2}(-z)^{5/2}} + \mathcal{O}(z^{-4}), & \text{if } \vartheta \neq 0. \end{cases}$$

Proof. These are proved using the asymptotics (3.8) of $\text{Ai}(t)$ and $\text{Bi}(t)$ as $t \rightarrow \infty$; for details, see Clarkson, Loureiro, and Van Assche [24]. \square

The plots in Figure 3.1 suggest the following conjecture (see also [24]).

Conjecture 3.4. If $q_n(z; 0)$ is defined by (3.7), then for $z < 0$ and $n \geq 1$, $q_n(z; 0)$ is a monotonically decreasing function and

$$q_{n+1}(z; 0) < q_n(z; 0).$$

3.2 Painlevé XXXIV equation and Jimbo-Miwa-Okamoto σ equation

Due to the one-to-one correspondence between solutions of P_{II} (1.2) and those of P_{34} (1.3) and S_{II} (1.4), as discussed in §2.1, we have the following result.

Theorem 3.5. *Equations (1.3) and (1.4) possess one-parameter family of solutions expressible in terms of Airy functions given by (3.4) if and only if $\alpha = n - \frac{1}{2}$, with $n \in \mathbb{Z}$.*

As for the Airy solutions of P_{II} (1.2), the Airy solutions of P_{34} (1.3) and S_{II} (1.4) can be expressed in terms of the determinant (3.6).

Theorem 3.6. *Let $\tau_n(z; \vartheta)$ be the determinant (3.6) for $n \geq 1$, with $\varphi(z; \vartheta)$ given by (3.4) and $\tau_0(z; \vartheta) = 1$, then*

$$p_n(z; \vartheta) = -2 \frac{d^2}{dz^2} \ln \tau_n(z; \vartheta), \quad (3.9)$$

$$\sigma_n(z; \vartheta) = \frac{d}{dz} \ln \tau_n(z; \vartheta), \quad (3.10)$$

satisfy P_{34} (1.3) and S_{II} (1.4) with $\alpha = n - \frac{1}{2}$, respectively.

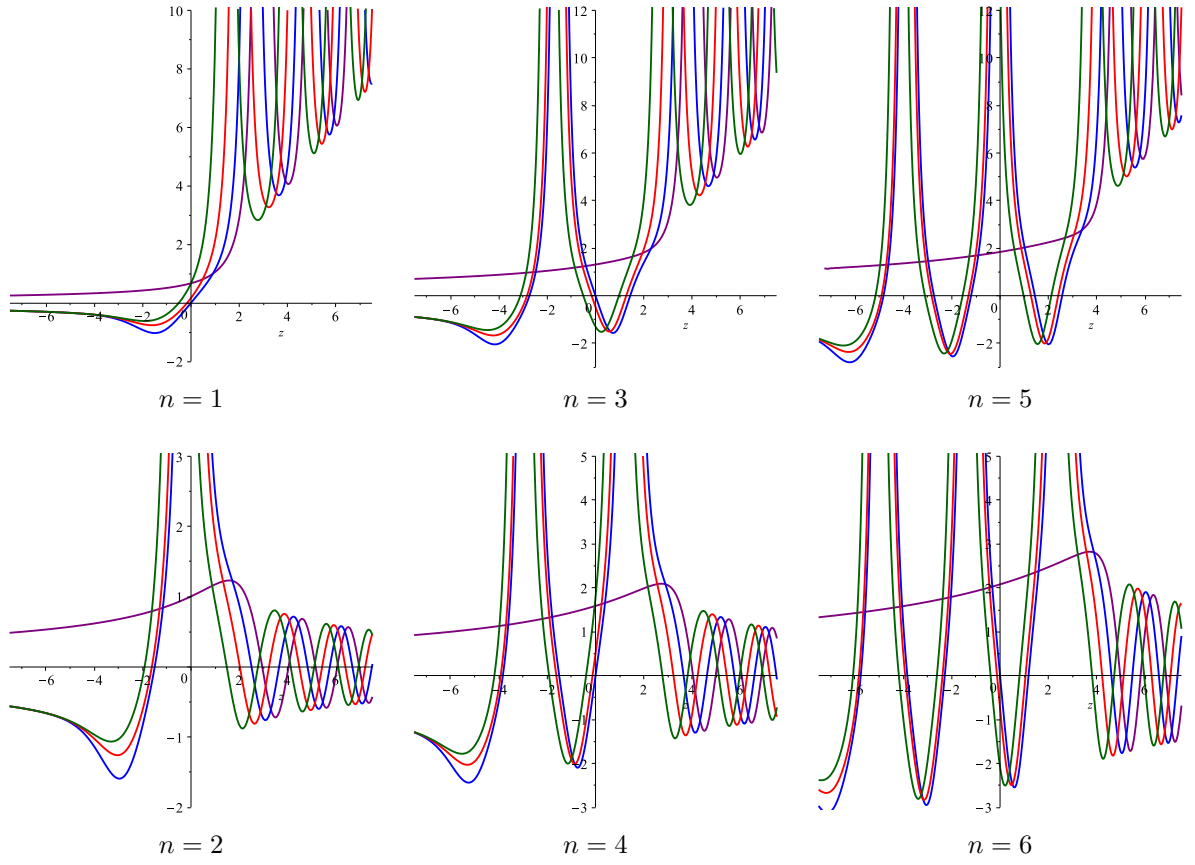


Figure 3.3: Plots of $p_n(z; \vartheta)/n$, with $p_n(z; \vartheta)$ given by (3.9), for $\vartheta = 0$ [purple], $\vartheta = \frac{1}{6}\pi$ [blue], $\vartheta = \frac{1}{3}\pi$ [red], $\vartheta = \frac{1}{2}\pi$ [green].

n	$\sigma_n(z; \vartheta)$
1	$-\Phi$
2	$\frac{1}{2\Phi^2 + z}$
3	$\frac{2z\Phi^2 - \Phi + z^2}{4\Phi^3 + 2z\Phi + 1}$
4	$\frac{48\Phi^3 - 8z^2\Phi^2 + 28z\Phi - 4z^3 + 9}{z(8z\Phi^4 - 16\Phi^3 + 8z^2\Phi^2 - 8z\Phi + 2z^3 - 3)} + \frac{3}{z}$

Table 3.3: The Airy solutions $\sigma_n(z; \vartheta)$ (3.10).

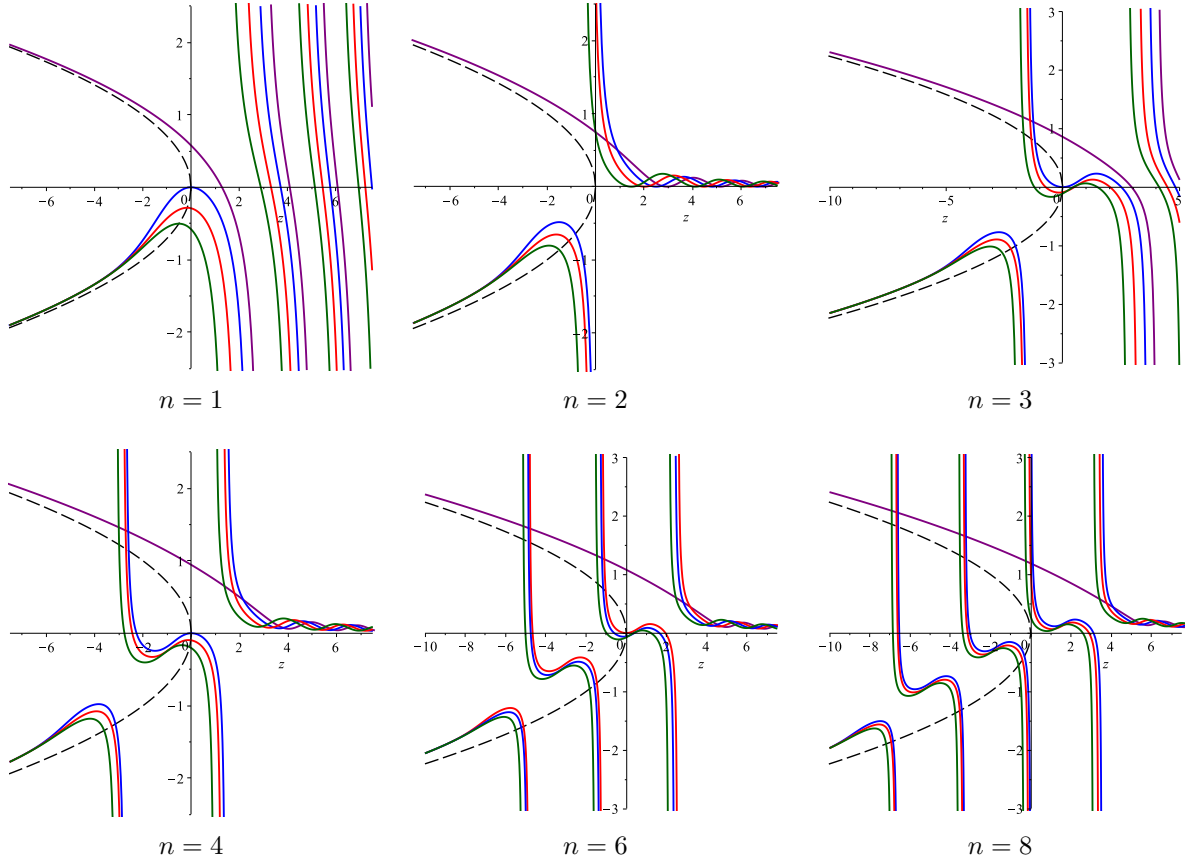


Figure 3.4: Plots of $\sigma_n(z; \vartheta)/n$, with $\sigma_n(z; \vartheta)$ given by (3.10), for $\vartheta = 0$ [purple], $\vartheta = \frac{1}{6}\pi$ [blue], $\vartheta = \frac{1}{3}\pi$ [red], $\vartheta = \frac{1}{2}\pi$ [green]; the dashed line is the parabola $2\sigma^2 + z = 0$.

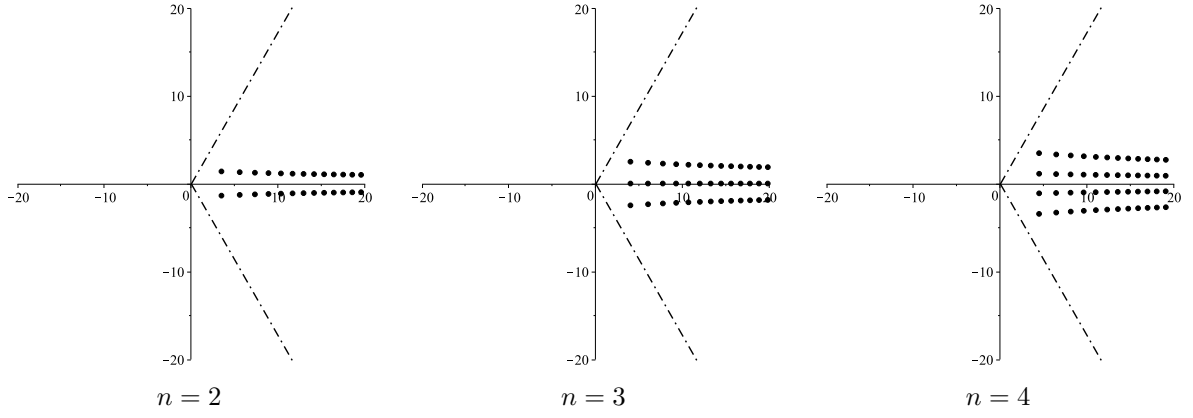


Figure 3.5: Plots of the poles of $p_n(z; 0)$ and $\sigma_n(z; 0)$ for $n = 2, 3, 4$.

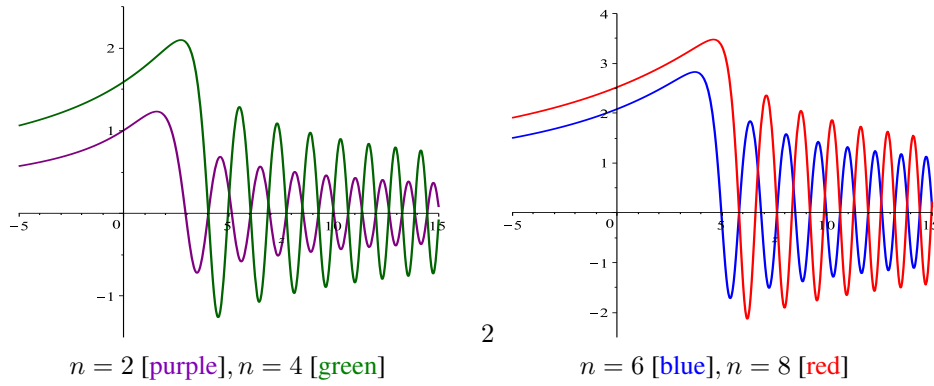


Figure 3.6: Plots of $p_n(z; 0)$ (3.9), for $n = 2, 4, 6, 8$.

Comparing (3.7) and (3.10), we see that

$$q_n(z; \vartheta) = \sigma_{n-1}(z; \vartheta) - \sigma_n(z; \vartheta). \quad (3.11)$$

The first few Airy solutions of P_{34} (1.3) and S_{II} (1.4) are given in Tables 3.2 and 3.3, respectively, and plots in Figures 3.3 and 3.4, respectively. As was the case for P_{II} , these plots show that the asymptotic behaviour as $z \rightarrow -\infty$ of the Airy solutions is completely different when $\vartheta = 0$ compared to the case when $\vartheta \neq 0$. The solutions $p_n(z; 0)$ and $\sigma_n(z; 0)$, which depends only on $\text{Ai}(t)$ and have the same poles, are *tronquée* solutions, see Figure 3.5. Further we see that the solutions $p_{2n}(z; 0)$ and $\sigma_{2n}(z; 0)$, for $n \in \mathbb{Z}$, are special in that they have no poles on the real axis, see also Figures 3.6 and 3.7. Additionally these solutions have oscillatory behaviour with algebraic decay as $z \rightarrow \infty$ as given in Theorem 3.8, see also Figures 3.8 and 3.9.

The asymptotic behaviour of the Airy solutions $p_n(z; \vartheta)$ and $\sigma_n(z; \vartheta)$ as $z \rightarrow -\infty$ is given in the following theorem.

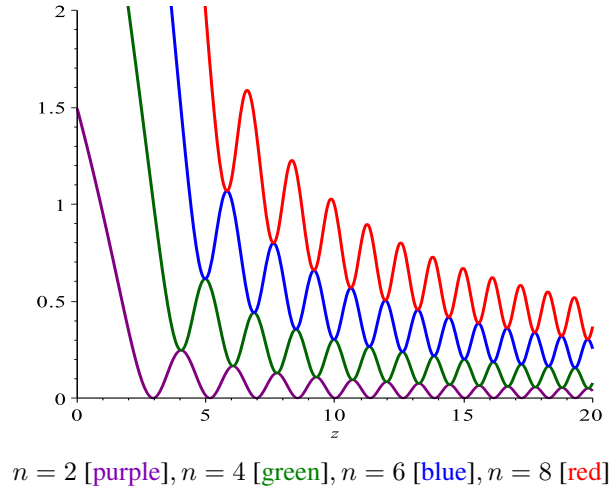


Figure 3.7: Plots of $\sigma_n(z; 0)$ (3.10), for $n = 2, 4, 6, 8$.

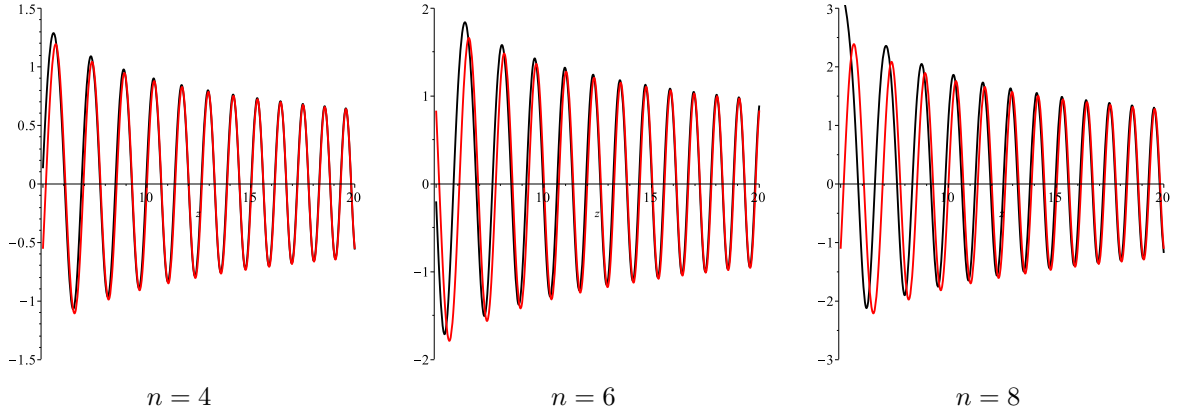


Figure 3.8: Plots of $p_n(z; 0)$ (3.9) [blue] and the leading term in the asymptotic expansion (3.14) [red] for $n = 4, 6, 8$.

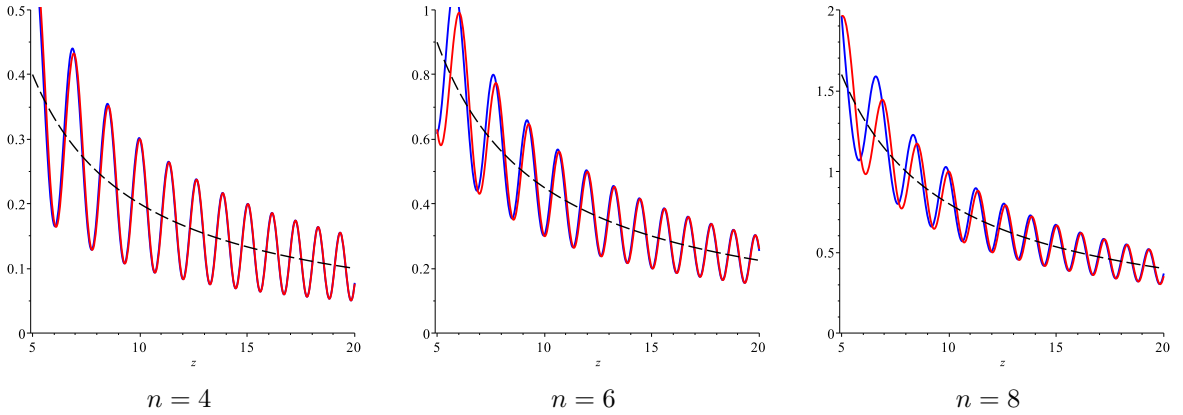


Figure 3.9: Plots of $\sigma_n(z; 0)$ (3.10) [blue] and the leading term in the asymptotic expansion (3.15) [red] for $n = 4, 6, 8$; the dashed line is $n^2/(8z)$.

Theorem 3.7. Let $p_n(z; \vartheta)$ be defined by (3.9) and $\sigma_n(z; \vartheta)$ by (3.10), then as $z \rightarrow -\infty$,

$$p_n(z; \vartheta) = \begin{cases} \frac{n}{\sqrt{2}(-z)^{1/2}} - \frac{n^2}{2z^2} + \frac{5n(4n^2+1)}{16\sqrt{2}(-z)^{7/2}} + \mathcal{O}(z^{-5}), & \text{if } \vartheta = 0, \\ -\frac{n}{\sqrt{2}(-z)^{1/2}} - \frac{n^2}{2z^2} - \frac{5n(4n^2+1)}{16\sqrt{2}(-z)^{7/2}} + \mathcal{O}(z^{-5}), & \text{if } \vartheta \neq 0, \end{cases} \quad (3.12)$$

$$\sigma_n(z; \vartheta) = \begin{cases} \frac{n(-z)^{1/2}}{\sqrt{2}} - \frac{n^2}{4z} - \frac{n(4n^2+1)}{16\sqrt{2}(-z)^{5/2}} + \mathcal{O}(z^{-4}), & \text{if } \vartheta = 0, \\ -\frac{n(-z)^{1/2}}{\sqrt{2}} - \frac{n^2}{4z} + \frac{n(4n^2+1)}{16\sqrt{2}(-z)^{5/2}} + \mathcal{O}(z^{-4}), & \text{if } \vartheta \neq 0. \end{cases} \quad (3.13)$$

Proof. The asymptotic expansions (3.12) and (3.13) are easily derived using the expansion for $q_n(z; \vartheta)$ given in Theorem 3.3 since if q is solution of P_{II} (1.2), then

$$p = \frac{dq}{dz} + q^2 + \frac{1}{2}z, \\ \sigma = \frac{1}{2} \left(\frac{dq}{dz} \right)^2 - \frac{1}{2}q^4 - \frac{1}{2}zq^2 - \left(\alpha + \frac{1}{2} \right)q - \frac{1}{8}z^2,$$

are solutions of P_{34} (1.3) and S_{II} (1.4), respectively \square

We remark that Its, Kuijlaars, and Östensson [25] noted that $p_2(z; 0)$ is special among all solutions $p_2(z; \vartheta)$ in its behavior as $z \rightarrow -\infty$.

The asymptotic behaviour of the solutions $p_n(z; \vartheta)$ and $\sigma_n(z; \vartheta)$ as $z \rightarrow \infty$, for $n \in 2\mathbb{Z}$, is given in the following theorem.

Theorem 3.8. Let $p_n(z; \vartheta)$ be defined by (3.9) and $\sigma_n(z; \vartheta)$ by (3.10) and $n \in 2\mathbb{Z}$, then as $z \rightarrow \infty$,

$$p_n(z; 0) = \frac{1}{2}n\sqrt{2}z^{-1/2} \cos\left(\frac{4}{3}\sqrt{2}z^{3/2} - \frac{1}{2}n\pi\right) + o(z^{-1/2}), \quad (3.14)$$

$$\sigma_n(z; 0) = \frac{n}{8z} \left\{ n - 2 \sin\left(\frac{4}{3}\sqrt{2}z^{3/2} - \frac{1}{2}n\pi\right) \right\} + o(z^{-1}). \quad (3.15)$$

Proof. Using the known asymptotic expansion of the Airy function as $t \rightarrow -\infty$, i.e.

$$\text{Ai}(t) = \pi^{-1/2}(-t)^{-1/4} \sin\left\{\frac{2}{3}(-t)^{3/2} + \frac{1}{4}\pi\right\} + o((-t)^{-1/4})$$

[23, §9.7(ii)] it is straightforward to verify the asymptotic expansions (3.14) and (3.15) for small values of n . \square

Plots of $p_n(z; 0)$ (3.9) and the leading term in the asymptotic expansion (3.14) for $n = 4, 6, 8$ are given in Figure 3.8. Analogous plots of $\sigma_n(z; 0)$ (3.10) and the leading term in the asymptotic expansion (3.15) for $n = 4, 6, 8$ are given in Figure 3.9.

Remark 3.9. Its, Kuijlaars, and Östensson [25] discussed solutions of the equation

$$u_\beta \frac{d^2 u_\beta}{dt^2} = \frac{1}{2} \left(\frac{du_\beta}{dt} \right)^2 + 4u_\beta^3 + 2tu_\beta^2 - 2\beta^2, \quad (3.16)$$

where β is a constant, which is equivalent to P_{34} (1.3) through the transformation $p(z) = 2^{1/3}u_\beta(t)$, with $t = -2^{-1/3}z$, and $\beta = \frac{1}{2}\alpha + \frac{1}{4}$ in their study of the double scaling limit of unitary random matrix ensembles of the form

$$Z_{n,N}^{-1} |\det M|^{2\beta} \exp\{-N \text{Tr} V(M)\} dM,$$

with $\beta > -\frac{1}{2}$ and V real analytic. In particular,

$$u_1(t) = -\frac{d^2}{dt^2} \mathcal{W}(\text{Ai}(t), \text{Ai}'(t)), \quad (3.17)$$

with $\mathcal{W}(\varphi_1, \varphi_2)$ the Wronskian, which is equivalent to the solution $p_2(z; 0)$ of P_{34} (1.3), and noted that on the positive real axis, $u_1(t)$ has an infinite number of zeros, which are the zeros of the Airy function $\text{Ai}(t)$, and an infinite number of additional zeros that interlace with the zeros of $\text{Ai}(t)$, see Figure 3.10.

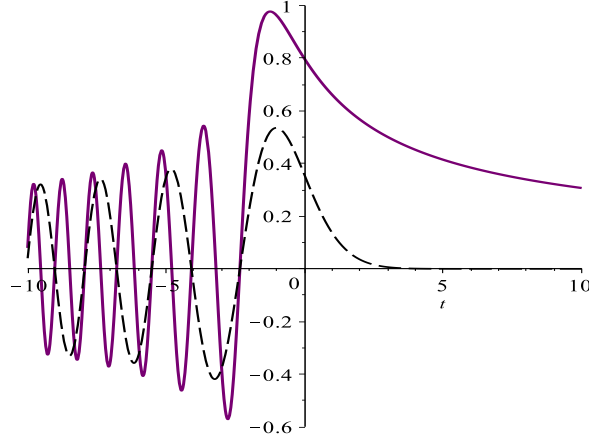


Figure 3.10: The solution $u_1(t)$ (3.17) of equation (3.16) [purple] and the Airy function $\text{Ai}(t)$ [dashed line].

Remark 3.10. Its, Kuijlaars, and Östensson [26, Theorem 1.2] prove that there are solutions $u_\beta(t)$ of (3.16) such that as $t \rightarrow \infty$

$$u_\beta(t) = \beta t^{-1/2} + \mathcal{O}(t^{-2}), \quad \text{as } t \rightarrow \infty. \quad (3.18)$$

$$u_\beta(t) = \beta(-t)^{-1/2} \cos \left\{ \frac{4}{3}(-t)^{3/2} - \beta\pi \right\} + \mathcal{O}(t^{-2}), \quad \text{as } t \rightarrow -\infty. \quad (3.19)$$

Letting $\beta = 1$ in (3.18) and (3.19) shows that they are in agreement with (3.12) and (3.14), for $n = 2$. Further Its, Kuijlaars, and Östensson conclude that solutions of (3.16) with asymptotic behaviour (3.18) are *tronquée* solutions, i.e. they have no poles in a sector of the complex plane.

The plots in Figures 3.3 and 3.4 suggest the following conjecture, which is analogous to Conjecture 3.4.

Conjecture 3.11. If $p_n(z; 0)$ is defined by (3.9) and $\sigma_n(z; 0)$ by (3.10), then for $z < 0$ and $n \geq 1$, $p_n(z; 0)$ is a monotonically decreasing function and $\sigma_n(z; 0)$ is a monotonically increasing function as z decreases. Further, for $z < 0$

$$p_{n+1}(z; 0) < p_n(z; 0), \quad \sigma_n(z; 0) < \sigma_{n+1}(z; 0).$$

4 Conclusion

In this paper we have studied the Airy solutions of P_{II} (1.2), P_{34} (1.3), and S_{II} (1.2). These show that when the solutions depend only on the Airy function $\text{Ai}(t)$, with $t = -2^{-1/3}z$, have a completely different asymptotic behaviour as $z \rightarrow -\infty$ compared to solutions which involve the Airy function $\text{Bi}(t)$. The special solutions which depend only on $\text{Ai}(t)$ are *tronquée* solutions, i.e. they have no poles in a sector of the complex plane. Further for both P_{34} (1.3) and S_{II} (1.2), it is shown that amongst these *tronquée* solutions there is a family of solutions which have no poles on the real axis. Such solutions of Painlevé equations often have important physical applications. For example, the well-known Hastings-McLeod solution of P_{II} (1.2) with $\alpha = 0$ [27], which arises in numerous applications, is a *tronquée* solution with no poles on the real axis, cf. [28, 29, 30].

The special solutions $q_n(z; 0)$ and $p_n(z; 0)$ of P_{II} and P_{34} arise in recent study by Clarkson, Loureiro, and Van Assche [24] of the discrete system

$$\begin{aligned} a_n + a_{n+1} &= b_n^2 - t, \\ a_n(b_n + b_{n-1}) &= n, \end{aligned}$$

which is equivalent to alt-dP_I (2.7) and arises in the study of orthogonal polynomials with respect to an exponential cubic weight [31, 32, 33].

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